# Coupled-mode theory for a plane and a cylindrical wave in optical waveguide Bragg grating lenses

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The coupling of a plane wave and a cylindrical wave, as expressed by a Hankel function, in optical waveguide Bragg grating lenses is theoretically treated in two dimensions. Analytical solutions of the coupled-mode equations are derived, and numerical examples for lenses with rectangular corrugations are shown for both the TE and the TM modes. The results are also compared with those obtained by treatment in one dimension.

#### INTRODUCTION

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A Bragg grating lens¹ for integrated optics is considered as a volume hologram that has the structure of an interference pattern produced by a plane and a cylindrical guided wave or by two cylindrical guided waves. The properties of volume holograms have been discussed by many authors. Kogelnik first treated the coupled-wave theory for infinite plane waves.<sup>3</sup> Solymsr extended the theory in two dimensions.<sup>3–3</sup> Kenan applied the theory to guided optical waves<sup>9</sup> and analyzed the operation of crossed-beam gratings made by two finite plane waves.<sup>19</sup>

In this paper, expressions for plane and cylindrical guided waves in a dielectric three-layer configuration are given, and the coupling of these waves in a Bragg grating lens, consisting of a rectangular corrugation in the waveguide film, is theoretically reasted in two dimensions. Solutions of the coupled-mode equations in integral forms are derived, and their numerical results are shown. A comparison of these results with those obtained in a one-dimensional analysis is also given.

# 1. WAVEGUIDE STRUCTURE AND GUIDED MODE

Consider a dielectric planar waveguide, as shown in Fig. 1, where d is the thickness of the film and  $n_1, n_2,$  and  $n_3$  are the refractive indices of the substrate, film, and cover, respectively; that is,

$$n(x) = \begin{cases} n_1 & x \leq 0, \\ n_2 & 0 \leq x \leq d, \\ n_3 & x \geq d, \end{cases}$$
(1)

and  $n_2 > n_5$ ,  $n_3$ . When time dependence  $\exp(i\omega t)$  of the wave field is assumed, Maxwell's equations in these media are written as follows:

$$\nabla \times \mathbf{H} = i\omega \epsilon_0 n^2(z) \mathbf{E}, \tag{2}$$

$$\nabla \times \mathbf{E} = i\omega \mu_0 \mathbf{H}$$
, (3)

$$\nabla \cdot \mathbf{E} = 0$$
, (4)

where E, H,  $\omega$ ,  $\epsilon_0$ , and  $\mu_0$  are the electric vector, the magnetic vector, the angular frequency, the dielectric constant of vacuum, and the magnetic permeability of vacuum, respectively. Here, the relative magnetic permeabilities of these media are assumed to be unity. Equations (2)–(4) lead to the wave equation for the electric field E, given by

$$\nabla^{2}\mathbf{E} + h_{0}^{2}n^{2}(z)\mathbf{E} = 0,$$
 (5)

where  $h_0^2 = \omega^2 \epsilon_0 \mu_0$ .

In this section, only the guided wave of the TE mode  $(E_s = 0)$  is considered. Then, a solution in the form

$$\mathbf{E} = \begin{pmatrix} A_x \\ A_y \\ 0 \end{pmatrix} F(x) \qquad (6)$$

is obtained, 11,12 where  $A_j(j=x,y)$  must satisfy the equations

$$\left(\frac{\partial^{2}}{\partial x^{2}} + \frac{\partial^{2}}{\partial y^{2}}\right)A_{j} + \beta^{2}A_{j} = 0, \quad (7)$$

$$\frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} = 0,$$
 (8)

and.

$$F(x) = \begin{cases} \cos \Phi_1 \exp(\gamma_1 x) & x \le 0, \\ \cos(\kappa_3 x - \Phi_1) & 0 \le x \le d, \\ (-1)^m \cos \Phi_3 \exp[-\gamma_3 (x - d)] & x \ge d, \end{cases}$$
(9)

$$k_0^2 n_1^2 + \gamma_1^2 = k_0^2 n_2^2 - \kappa_2^2 = k_0^2 n_3^2 + \gamma_3^2 = \beta^2$$
, (10)

$$\kappa_2 d - \Phi_1 - \Phi_3 = m\pi(m, integer),$$
 (11)

$$\tan \Phi_1 = \frac{\gamma_1}{\kappa_2}$$
,  $\tan \Phi_3 = \frac{\gamma_3}{\kappa_2}$  (12)

For a plane wave propagating in the  $\theta$  direction from the x axis, the electric and magnetic fields, respectively, are given as

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